

Motivation

APPLICATIONS

Hydrogen Propulsion:

Explosion risk assesment

Chemical Manufacturing:

Reaction control

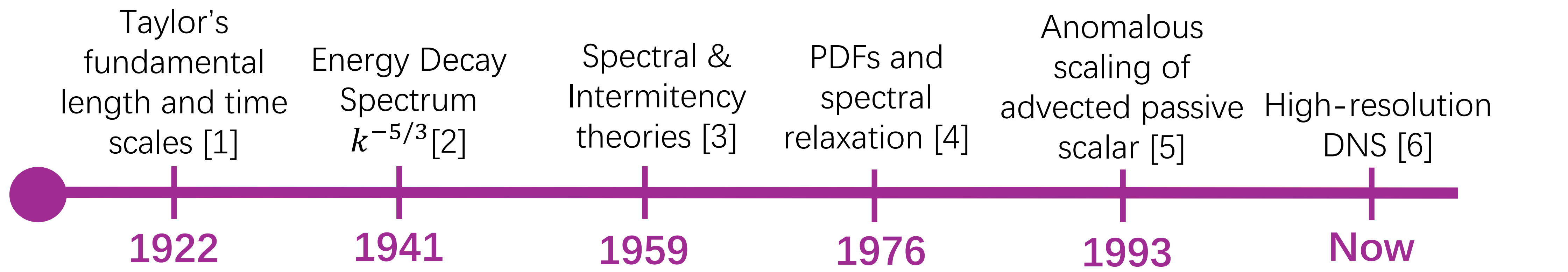
Environmental Protection:

Toxic exposure limits

Mixing in turbulent flows is notoriously **difficult to predict** due to extreme fluctuations and highly variable mixing rates. Rather than focusing on averages, this research models the **maximum of the scalar**.

The advantage is that by defining this upper limit, we can determine if a safety threshold will be crossed; if the critical limit is above the predicted maximum, the probability of an incident is mathematically zero.

State of the Art



LIMITATIONS OF CURRENT MODELS

Statistical models

fail to capture transient extremes.

DNS give precise statistics but are

computationally very expensive

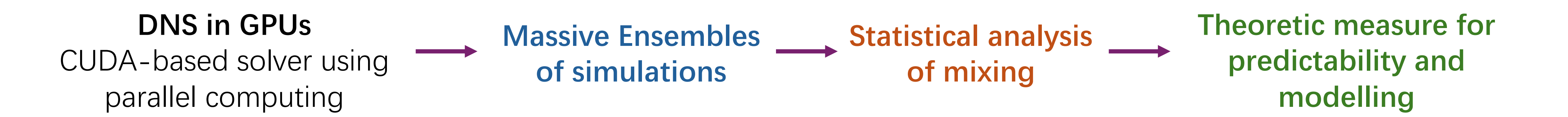
Most studies focus on steady-state mixing and

not transient

Objectives

- Construct an extensive **database** describing the statistics of the transient mixing of scalar in turbulence
- Characterization of transient **mixing dynamics** in terms of flow patterns
- Construct **statistical models** that provide theoretical explanations and predict statistics of transient mixing depending on the initial distribution of scalar

Methodology



EQUATIONS

Mass conservation :

$\nabla \cdot \mathbf{u} = 0$

Momentum conservation:

$\frac{\partial \mathbf{u}}{\partial t} + \boldsymbol{\omega} \times \mathbf{u} = -\frac{1}{\rho} \nabla p + \nu \nabla^2 \mathbf{u} + \mathbf{c} \mathbf{f}$

Fick's law:

$\frac{\partial \phi}{\partial t} + \mathbf{u} \cdot \nabla \phi = D \nabla^2 \phi$

$\mathbf{u} \equiv$ velocity vector

$\boldsymbol{\omega} \equiv$ vorticity vector

$p \equiv$ pressure

$\rho \equiv$ density

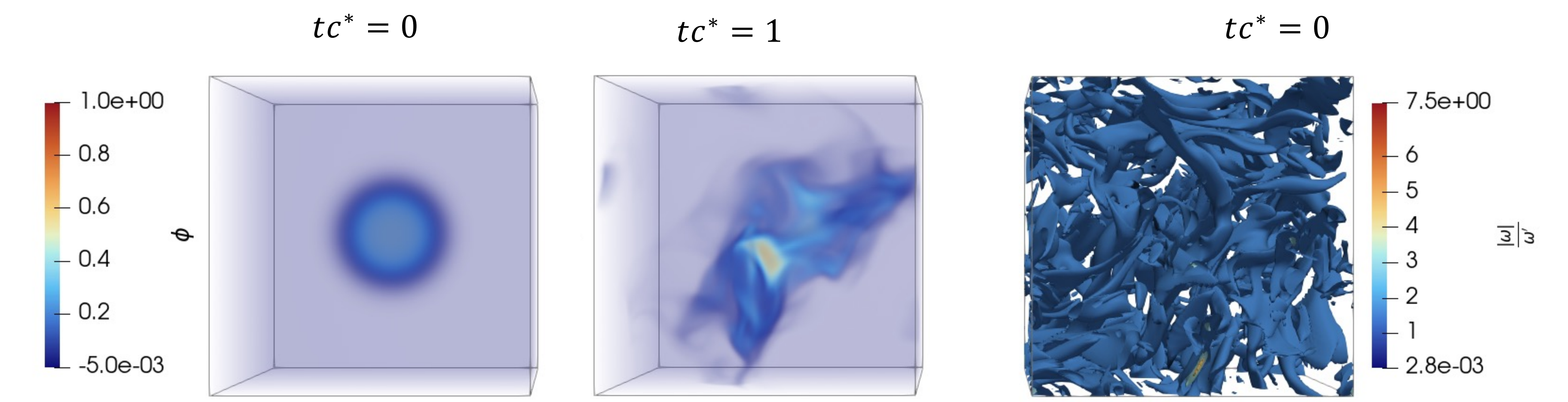
$\nu \equiv$ kinematic viscosity

$\mathbf{c} \mathbf{f} \equiv$ adaptive force to sustain turbulence

$\phi \equiv$ scalar mass fraction

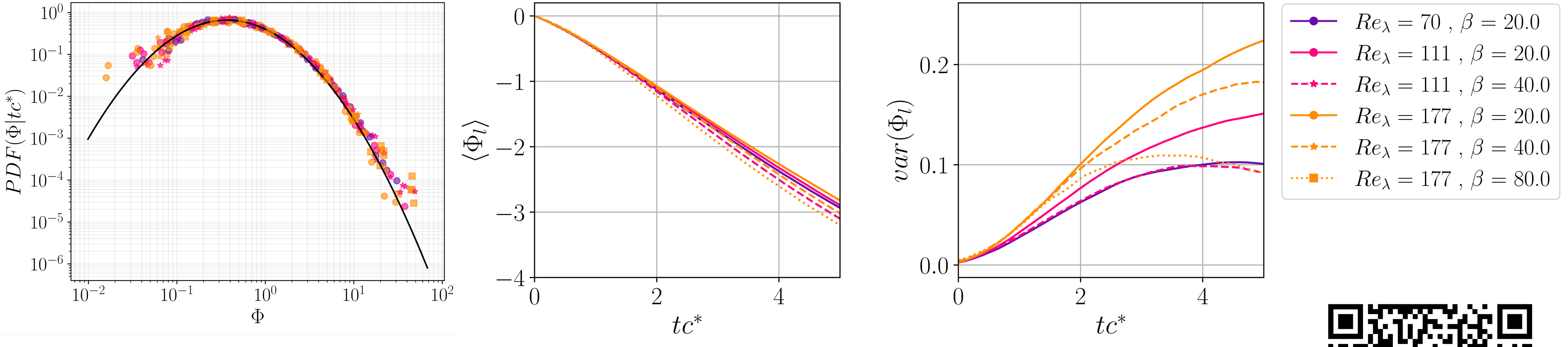
$D \equiv$ scalar diffusivity

- Initial condition of scalar: **Gaussian sphere of radius r_0**
- Mixing depends on Re_λ & $\beta = r_0/\eta$ (initial scalar size in Kolmogorov units)



- To account for the scalar size, time is non-dimensionalised with the inertial scalar time and a correction τ_0 is substracted to account for the time required to break the Gaussian : $tc^* = \frac{t-\tau_0}{t_i}$

Preliminary Research/Results



- Probability Density Function of $\Phi = \max(\phi)$ standarized at $tc^*[1.5, 2, 2.5, 3, 3.5, 4]$ follow a lognormal distribution $\longrightarrow$ Analyse statistics of $\Phi_l = Ln(\Phi)$
- The lognormal median shows an excellent collapse across all cases with scaling behaviour suggesting the presence of intermittency.
 - The lognormal variance fails to collapse $\rightarrow$ dependence on Re_λ and β



SCALAR MIXING & TURBULENT FLOW VIDEOS:

[1] Taylor, G. I. (1922). Diffusion by continuous movements. Proc. London Math. Soc., 2(1), 196-212.
[2] Andrei Nikolaevich Kolmogorov, V. Levin, Julian Charles Roland Hunt, Owen Martin Phillips, David Williams; Dissipation of energy in the locally isotropic turbulence. Proc. A 1 July 1991; 434 (1890): 15-17.
[3] Batchelor, G. K. (1959). Small-scale variation of convected quantities like temperature in turbulent fluid Part 1. General discussion and the case of small cond. J. Fluid Mech., 5(1), 113-133. (4), 437-450.
[4] Dopazo, C. (1976). A probabilistic approach to turbulent flame theory. Acta Astronautica, 3(9-10), 853-878.
[5] Kraichnan, R. H. (1994). Anomalous scaling of a randomly advected passive scalar. Phys. Rev. Lett., 72(7), 1016-1019.
[6] Schumacher, J., Sreenivasan, K. R., & Yeung, P. K. (2005). Very fine structures in scalar mixing. Journal of Fluid Mechanics, 531, 113-122.